

Renormalization group theory and its application to highly oscillatory problems

Wenlei Li

School of mathematics, JLU. China

Joint work with: Shaoyun Shi (JLU)

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Introduction & motivations.



Renormalization group (RG) theory

- [Chen *et al* \(Phys. Rev. Lett. 1994\)](#) Renormalization group theory for global asymptotic analysis.
 - [Chen *et al* \(Phys. Rev. E. 1996\)](#) Renormalization group and singular perturbations: multiple scales, boundary layers, and reductive perturbation theory.
- **Main goal:** To compute the **effective approximate solution** of different kinds of singular perturbation problems **in a unified manner**.
 - **Advantage:** It starts with a naive expansion and **does not require any further a priori assumptions regarding the structure of the perturbation series**, like an anticipation of scales involved, as is done in multiple scale analysis



RG theory: recent works

- [Ei *et al* \(Ann. Physics, 2000\)](#) Renormalization-group method for reduction of evolution equations: invariant manifolds and envelopes.
- [Ziane \(J. Math. Phys 2003\)](#) On a certain renormalization group method.
- [Goldenfeld *et al* \(J. Statis. Phys. 2006\)](#) Renormalization Group Approach to Multiscale Modelling in Materials Science.
- [Chiba \(SIAM J. Appl. Dyn. Syst. 2009\)](#) Extension and Unification of Singular Perturbation Methods for ODEs Based on the Renormalization Group Method.
- [Kirkinis \(SIAM Review. 2012\)](#) The Renormalization Group: A Perturbation Method for the Graduate Curriculum.
- ...



RG theory: two cases

Ziane, 2003, JMP

$$\begin{cases} \dot{\mathbf{x}} + \frac{1}{\varepsilon} A \mathbf{x} = \mathbf{f}(\mathbf{x}), \\ \mathbf{x}(0) = \mathbf{x}_0, \end{cases}$$

- A : a complex diagonalizable matrix;
- $\mathbf{f}(\mathbf{x})$: a polynomial nonlinear term.

Chiba, 2009, SIAM. JADS

$$\dot{\mathbf{x}} = \varepsilon \mathbf{g}(\mathbf{x}, t, \varepsilon), \quad \mathbf{x} \in U,$$

- U : an open set in \mathbb{C}^n , the closure \bar{U} is compact;
- $\mathbf{g}(\mathbf{x}, t, \varepsilon)$: almost-periodic in t , and the set of corresponding Fourier exponents having no accumulation on \mathbb{R} .

Motivation 1

- 1 How about the theory for **more general** cases?
- 2 How about the theory for **other particular** cases?



Highly oscillatory second order ODEs

$$\begin{cases} \varepsilon^2 \ddot{\mathbf{y}}(t) + (A + \frac{1}{\varepsilon^2} B) \mathbf{y}(t) + \mathbf{f}(\mathbf{y}(t)) = 0, \mathbf{y} \in \mathbb{C}^d, t > 0, \\ \mathbf{y}(0) = \eta_1, \dot{\mathbf{y}}(0) = \frac{\eta_2}{\varepsilon^2}, \end{cases} \quad (1)$$

- ε : small real parameter $|\varepsilon| \ll 1$;
- $A(B)$: nonnegative (positive) definite matrix;
- $\mathbf{f}(e^{is} \mathbf{y}) = e^{is} \mathbf{f}(\mathbf{y}), \forall s \in \mathbb{R}$.

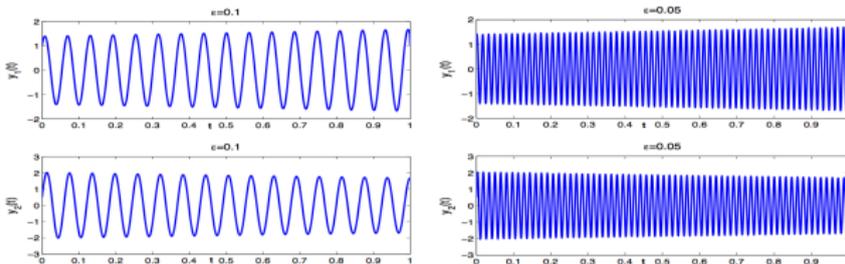
Nonlinear Klein-Gordon equation in the nonrelativistic limit regime

$$\begin{cases} \varepsilon^2 \partial_{tt} u - \Delta u + \frac{1}{\varepsilon^2} u + f(u) = 0, & \mathbf{x} \in \mathbb{R}^n, t > 0, \\ u(\mathbf{x}, 0) = \phi(\mathbf{x}), \quad \partial_t u(\mathbf{x}, 0) = \frac{1}{\varepsilon^2} \gamma(\mathbf{x}), \end{cases} \quad (2)$$

- ε : small real parameter $|\varepsilon| \ll 1$ inversely proportional to the speed of light;
- ϕ, γ : real-valued functions;
- $f(e^{is}u) = e^{is}f(u), \forall s \in \mathbb{R}$.

High oscillation property

Propagates highly oscillatory waves with wavelength at $O(\varepsilon^2)$ and amplitude at $O(1)$.



- Difficulties in the asymptotic analysis.
- Burdens in practical computation, making the approximation extremely challenging and costly in the regime of $0 < \varepsilon \ll 1$.

Known works

- [Machihara *et al.* \(Math. Ann. 2002\)](#) Nonrelativistic limit in the energy space for non-linear Klein-Gordon equations.
- [Masmoudi & Nakanishi \(Math. Ann. 2002\)](#) From nonlinear Klein-Gordon equation to a system of coupled nonlinear Schrödinger equations.
- [Bao & Dong \(Numer. Math. 2012\)](#) Analysis and comparison of numerical methods for the Klein-Gordon equation in the nonrelativistic limit regime.
- [Bao, Dong & Zhao \(J. Math. Study. 2014\)](#) Uniformly Accurate Multiscale Time Integrators for Highly Oscillatory Second Order Differential Equations.
- ...



Motivation 2.

- 1 To make a complete investigation of the asymptotic properties as $\varepsilon \rightarrow 0$, from the point of renormalization group theory.
- 2 To present a rigorous statement of the following assumption proposed to secure the validity of the numerical results:

$$\|\mathbf{y}^{(k)}(t)\| \leq \frac{M}{\varepsilon^{2k}}, \quad k = 0, 1, 2, \dots \quad x \in [0, T] \subset [0, T^*). \quad (3)$$

RG theory revisited.



KBM condition

$$\dot{\mathbf{x}} = \varepsilon \mathbf{g}(\mathbf{x}, t, \varepsilon), \quad \mathbf{x} \in U, \quad (4)$$

- U : open set in \mathbb{C}^n , and the closure \bar{U} is compact.

(G₁): $\mathbf{g}(\mathbf{x}, t, \varepsilon)$ is C^2 in $\mathbf{x} \in U$, C^1 in t and analytic in $\varepsilon \in I_0$, with $I_0 \subset \mathbb{R}$ an open neighborhood of the origin. Furthermore, $\mathbf{g}(\mathbf{x}, t, \varepsilon)$ is Lipschitz continuous in \mathbf{x} on U , i.e., there exists a constant $L_U > 0$ such that

$$\|\mathbf{g}(\mathbf{x}, t, \varepsilon) - \mathbf{g}(\mathbf{y}, t, \varepsilon)\| \leq L_U \|\mathbf{x} - \mathbf{y}\|,$$

for all $\mathbf{x}, \mathbf{y} \in U, t \in \mathbb{R}, \varepsilon \in I_0$.

KBM condition:

$$R(\mathbf{y}, \varepsilon) = \lim_{t \rightarrow \infty} \frac{1}{t - t_0} \int_{t_0}^t \mathbf{g}(\mathbf{y}, s, \varepsilon) ds \quad (5)$$

converges for arbitrary $t_0 \in \mathbb{R}_+$, and is uniform in $\mathbf{y} \in \bar{U}, \varepsilon \in I_0$.

RG step 1: Naive expansion

Expand \mathbf{x} as $\mathbf{x} = \mathbf{x}^{(0)} + \varepsilon \mathbf{x}^{(1)} + \varepsilon^2 \mathbf{x}^{(2)} + \dots$,

$$\dot{\mathbf{x}}^{(0)} = 0, \quad (6)$$

$$\dot{\mathbf{x}}^{(1)} = G_1(\mathbf{x}^{(0)}, t), \quad (7)$$

$$\dot{\mathbf{x}}^{(2)} = G_2(\mathbf{x}^{(0)}, \mathbf{x}^{(1)}, t), \quad (8)$$

\dots ,

with

$$G_1(\mathbf{x}^{(0)}, t) = \mathbf{g}(\mathbf{x}^{(0)}, t, 0),$$

$$G_2(\mathbf{x}^{(0)}, \mathbf{x}^{(1)}, t) = \frac{\partial \mathbf{g}}{\partial \mathbf{x}}(\mathbf{x}^{(0)}, t, 0) \mathbf{x}^{(1)} + \frac{\partial \mathbf{g}}{\partial \varepsilon}(\mathbf{x}^{(0)}, t, 0),$$

\dots



RG step 1: Naive expansion

$$\mathbf{x}^{(0)} = \xi_0$$

$$\mathbf{x}^{(1)} = \int_0^t G_1(\xi_0, s) ds = R_1(\xi_0)t + N_1(\xi_0, t),$$

$$\begin{aligned} \mathbf{x}^{(2)} &= \int_0^t G_2(\xi_0, R_1(\xi_0)s + N_1(\xi_0, s), s) ds \\ &= (R_2(\xi_0) + \frac{\partial N_1}{\partial \xi_0}(\xi_0, t)R_1(\xi_0))t + N_2(\xi_0, t) + \frac{1}{2} \frac{\partial R_1}{\partial \xi_0}(\xi_0)R_1(\xi_0)t^2, \end{aligned}$$

$$\vdots$$

with

$$R_1(\xi_0) = \lim_{t \rightarrow \infty} \frac{1}{t} \int_0^t G_1(\xi_0, s) ds,$$

$$N_1(\xi_0, t) = \int_0^t (G_1(\xi_0, s) - R_1(\xi_0)) ds;$$



RG step 1: Naive expansion

$$\begin{aligned}
 R_2(\xi_0) &= \lim_{t \rightarrow \infty} \frac{1}{t} \int_0^t (G_2(\xi_0, N_1(\xi_0, s), s) - \frac{\partial N_1}{\partial \xi_0}(\xi_0, s)R_1(\xi_0)) ds, \\
 N_2(\xi_0, t) &= \int_0^t (G_2(\xi_0, N_1(\xi_0, s), s) - \frac{\partial N_1}{\partial \xi_0}(\xi_0, s)R_1(\xi_0) - R_2(\xi_0)) ds, \\
 &\vdots
 \end{aligned}$$

$$\begin{aligned}
 \mathbf{x}(\xi_0, t, \varepsilon) &= \xi_0 + \varepsilon(R_1(\xi_0)t + N_1(\xi_0, t)) + \varepsilon^2((R_2(\xi_0) + \frac{\partial N_1}{\partial \xi_0}(\xi_0, t)R_1(\xi_0))t \\
 &\quad + N_2(\xi_0, t) + \frac{1}{2} \frac{\partial R_1}{\partial \xi_0}(\xi_0)R_1(\xi_0)t^2) + O(\varepsilon^3). \tag{9}
 \end{aligned}$$

RG step 2: Renormalization

Propose an ansatz free parameter $\mu \in \mathbb{R}$ into $\mathbf{x}(\xi_0, t, \varepsilon)$ to replace t^i by $t^i - \mu^i + \mu^i, i = 1, 2, \dots$.

$$\begin{aligned} \mathbf{x}(\xi_0, t, \varepsilon) = & \xi_0 + \varepsilon(R_1(\xi_0)(t - \mu + \mu) + N_1(\xi_0, t)) + \\ & \varepsilon^2((R_2(\xi_0) + \frac{\partial N_1}{\partial \xi_0}(\xi_0, t)R_1(\xi_0))(t - \mu + \mu) + \\ & N_2(\xi_0, t) + \frac{1}{2} \frac{\partial R_1}{\partial \xi_0}(\xi_0)R_1(\xi_0))(t^2 - \mu^2 + \mu^2)) + \dots \end{aligned}$$



RG step 2: Renormalization

Construct

$$\xi(\mu, \varepsilon) = \xi_0 + \varepsilon A_1(\mu) + \varepsilon^2 A_2(\mu) + \dots$$

such that

$$\mathbf{x}(\xi_0, t, \varepsilon) = \tilde{\mathbf{x}}(\xi(\mu, \varepsilon), t - \mu, t, \varepsilon), \quad (10)$$

without the secular terms and the ones as
 $\varepsilon R_1(\xi)\mu, \varepsilon^2((R_2(\xi) + \frac{\partial N_1}{\partial \xi}(\xi, t)R_1(\xi)))\mu, \dots$.



RG step 2: Renormalization

$$\xi_0 = \xi(\mu, \varepsilon) - \varepsilon R_1(\xi(\mu, \varepsilon))\mu - \varepsilon^2 (R_2(\xi(\mu, \varepsilon))\mu + \frac{\partial R_1}{\partial \xi}(\xi(\mu, \varepsilon))R_1(\xi(\mu, \varepsilon))\mu^2) + \dots,$$

$$\begin{aligned} \tilde{\mathbf{x}}(\xi(\mu, \varepsilon), t - \mu, t, \varepsilon) &= \mathbf{x}(\xi_0, t, \varepsilon) \\ &= \xi(\mu, \varepsilon) + \varepsilon (R_1(\xi(\mu, \varepsilon))(t - \mu) + N_1(\xi(\mu, \varepsilon), t)) + \\ &\quad \varepsilon^2 ((R_2(\xi(\mu, \varepsilon)) + \frac{\partial N_1}{\partial \xi}(\xi(\mu, \varepsilon), t)R_1(\xi(\mu, \varepsilon)))(t - \mu) \\ &\quad + N_2(\xi(\mu, \varepsilon), t) + \frac{1}{2} \frac{\partial R_1}{\partial \xi}(\xi(\mu, \varepsilon))R_1(\xi(\mu, \varepsilon))(t - \mu)^2) + \dots \end{aligned}$$

RG step 2: Renormalization

$$\frac{\partial}{\partial \mu} \tilde{\mathbf{x}}(\xi(\mu, \varepsilon), t - \mu, t, \varepsilon) = \frac{\partial}{\partial \mu} \mathbf{x}(\xi_0, t, \varepsilon) \equiv 0.$$

Especially, by taking $\mu = t$, we get the **renormalized equation**:

$$\dot{\xi} = \varepsilon R_1(\xi) + \varepsilon^2 R_2(\xi) + \dots \quad (11)$$

with the initial value condition $\xi(0) = \xi_0$, and the corresponding **approximate solution**:

$$\mathbf{x}(\xi_0, t, \varepsilon) = \xi(t) + \varepsilon N_1(\xi(t), t) + \varepsilon^2 N_2(\xi(t), t) + \dots$$



RG step 2: Asymptotic results

(First order) RG equation

$$\dot{\xi} = \varepsilon R_1(\xi). \quad (12)$$

(First order) approximate solution

$$\hat{\mathbf{x}}(t) = \xi(t) + \varepsilon N_1(\xi(t), t).$$



RG step 3: Estimation

By $\tau = \varepsilon t$, (12) becomes

$$\frac{d\eta}{d\tau} = R_1(\eta). \quad (13)$$

Theorem 1.

Assume that $\mathbf{g}(\mathbf{x}, t, \varepsilon)$ satisfies (G_1) and the KBM condition. Let $\mathbf{x}(t)$ be a solution of (4), $\eta(\tau)$ be the solution of (13) with $\eta(0) = \mathbf{x}(0)$ and the maximum existence interval (a, b) . Then, for any closed interval $[T_1, T_2] \subset (a, b)$, there exists a constant $\varepsilon_0 > 0$, such that, for any $\varepsilon \in (0, \varepsilon_0)$,

$$\|\mathbf{x}(t) - \hat{\mathbf{x}}(t)\| < C\delta(\varepsilon),$$

as long as $\frac{T_1}{\varepsilon} < t < \frac{T_2}{\varepsilon}$, where C is a positive constant independent to ε , and $\delta(\varepsilon) = o(1)$ as $\varepsilon \rightarrow 0$.

Periodic case

Corollary 1.

Assume that $\mathbf{g}(\mathbf{x}, t, \varepsilon)$ satisfies (G_1) , and it is periodic in t . Let $\mathbf{x}(t)$ be a solution of (4), $\eta(\tau)$ be the solution of (13) with $\eta(0) = \mathbf{x}(0)$ and the maximum existence interval (a, b) . Then, for any closed interval $[T_1, T_2] \subset (a, b)$, there exists a constant $\varepsilon_0 > 0$, such that for any $\varepsilon \in (0, \varepsilon_0)$,

$$\|\mathbf{x}(t) - \hat{\mathbf{x}}(t)\| < C\varepsilon, \quad (14)$$

as long as $\frac{T_1}{\varepsilon} < t < \frac{T_2}{\varepsilon}$, where C is a positive constant independent to ε .

- A same result can be found in [Chiba. SIAM J. Appl. Dyn. Syst. 2009].



Quasi-periodic case

(G₂): $\mathbf{g}(\mathbf{x}, t, \varepsilon)$ is sufficiently smooth in $(\mathbf{x}, t, \varepsilon)$ and quasi-periodic in t , i.e., it can be rewritten as the following form

$$\mathbf{g}(\mathbf{x}, t, \varepsilon) = \sum_{\mathbf{v} \in \mathbb{Z}^d} \mathbf{g}_{\mathbf{v}}(\mathbf{x}, \varepsilon) e^{i\langle \omega, \mathbf{v} \rangle t} = \sum_{\mathbf{v} \in \mathbb{Z}^d} \sum_{l=0}^{\infty} \varepsilon^l \mathbf{g}_{\mathbf{v}l}(\mathbf{x}) e^{i\langle \omega, \mathbf{v} \rangle t}, \quad (15)$$

where $\omega \in \mathbb{R}^d$ is a constant vector, $i = \sqrt{-1}$, and

$$\max\{\|\mathbf{g}_{\mathbf{v}l}(\mathbf{x})\|, \|\nabla \mathbf{g}_{\mathbf{v}l}(\mathbf{x})\|\} \leq \Gamma_0 \Gamma_1^l e^{-\sigma|\mathbf{v}|} \quad (16)$$

is valid for some $\Gamma_0, \Gamma_1, \sigma > 0$ and any $l \in \mathbb{N}$, $\mathbf{x} \in U$.

Diophantine conditions

Bryuno-Diophantine condition

For fixed vector $\eta \in \mathbb{R}^l$, let

$$\alpha_m^l(\eta) = \inf\{|\langle \eta, \nu \rangle| \mid \nu \in \mathbb{Z}^l, \text{ such that } 0 < |\nu| < 2^m\}, m \in \mathbb{N}.$$

Then the Bryuno function

$$\mathfrak{B}(\eta) = \sum_{m=0}^{\infty} \frac{1}{2^m} \ln \frac{1}{\alpha_m^l(\eta)} < \infty.$$

GBD condition

For fixed vector $\eta \in \mathbb{R}^l$, decompose η as $\eta = (a_1 \eta^{(1)}, \dots, a_j \eta^{(j)}) \in \mathbb{R}^{l_1} \times \dots \times \mathbb{R}^{l_j}$ with $\eta^i \in \mathbb{Q}^{l_i}, l_i \geq 0, i = 1, \dots, j, l_1 + \dots + l_j = l$, such that $\mathbf{a} = (a_1, \dots, a_j)$ satisfies the Bryuno - Diophantine condition.

Quasi-periodic case

Theorem 2.

Assume the assumption (G_2) holds, and the corresponding ω satisfies the GBD condition. Let $\mathbf{x}(t)$ be a solution of (4), $\eta(\tau)$ be the solution of (13) with $\eta(0) = \mathbf{x}(0)$ and the maximum existence interval (a, b) . Then, for any closed interval $[T_1, T_2] \subset (a, b)$, there exists a constant $\varepsilon_0 > 0$, such that for any $\varepsilon \in (0, \varepsilon_0)$,

$$\|\mathbf{x}(t) - \hat{\mathbf{x}}(t)\| < C\varepsilon, \quad (17)$$

as long as $\frac{T_1}{\varepsilon} < t < \frac{T_2}{\varepsilon}$, where C is a positive constant independent to ε .



Highly oscillatory problems.



Model

$$\begin{cases} \varepsilon^2 \ddot{\mathbf{y}}(t) + (A + \frac{1}{\varepsilon^2} B) \mathbf{y}(t) + \mathbf{f}(\mathbf{y}(t)) = 0, \mathbf{y} \in \mathbb{C}^d, t > 0, \\ \mathbf{y}(0) = \eta_1, \dot{\mathbf{y}}(0) = \frac{\eta_2}{\varepsilon^2}, \end{cases} \quad (18)$$

- ε : small real parameter $|\varepsilon| \ll 1$;
 - $A(B)$: nonnegative (positive) definite matrix;
 - $\mathbf{f}(e^{is}\mathbf{y}) = e^{is}\mathbf{f}(\mathbf{y}), \forall s \in \mathbb{R}$.
-
- Without loss of generality, we assume $B = \Lambda^2$ with $\Lambda = \text{diag}\{\lambda_1, \dots, \lambda_d\}, \lambda_j > 0, j = 1, \dots, d..$



Transformed problems

$$\text{By } \epsilon = \varepsilon^2, \quad t = \epsilon x, \quad \mathbf{y}_1(x) = \mathbf{y}(t), \quad \mathbf{y}_2(x) = \mathbf{y}'_1(x), \quad (19)$$

where ' denotes the derivation with respect to x .

$$\begin{cases} \begin{pmatrix} \mathbf{y}'_1 \\ \mathbf{y}'_2 \end{pmatrix} = \begin{pmatrix} O & E \\ -B & O \end{pmatrix} \begin{pmatrix} \mathbf{y}_1 \\ \mathbf{y}_2 \end{pmatrix} + \epsilon \begin{pmatrix} O \\ -A\mathbf{y}_1 - \mathbf{f}(\mathbf{y}_1) \end{pmatrix}, \\ \begin{pmatrix} \mathbf{y}_1(0) \\ \mathbf{y}_2(0) \end{pmatrix} = \begin{pmatrix} \eta_1 \\ \eta_2 \end{pmatrix}, \end{cases} \quad (20)$$

Transformed problems

$$\text{By } \begin{pmatrix} \mathbf{y}_1 \\ \mathbf{y}_2 \end{pmatrix} = \begin{pmatrix} E & i\Lambda \\ iE & \Lambda \end{pmatrix} \begin{pmatrix} e^{i\Lambda x} & O \\ O & e^{-i\Lambda x} \end{pmatrix} \begin{pmatrix} \mathbf{z}_1 \\ \mathbf{z}_2 \end{pmatrix}. \quad (21)$$

Then (20) becomes

$$\begin{cases} \mathbf{z}' = \epsilon \mathbf{g}(\mathbf{z}, x), \\ \mathbf{z}(0) = \begin{pmatrix} \frac{\eta_1 - i\Lambda^{-1}\eta_2}{2} \\ \frac{-i\eta_1 + \Lambda^{-1}\eta_2}{2} \end{pmatrix}. \end{cases} \quad (22)$$

$$\mathbf{g}(\mathbf{z}, x) = \frac{1}{2} \begin{pmatrix} A(i\Lambda^{-1}\mathbf{z}_1 - e^{-2i\Lambda x}\mathbf{z}_2) + i\Lambda^{-1}e^{-i\Lambda x}\mathbf{f}(e^{i\Lambda x}\mathbf{z}_1 + i\Lambda e^{-i\Lambda x}\mathbf{z}_2) \\ -A(\Lambda^{-1}e^{2i\Lambda x}\mathbf{z}_1 + i\mathbf{z}_2) - \Lambda^{-1}e^{i\Lambda x}\mathbf{f}(e^{i\Lambda x}\mathbf{z}_1 + i\Lambda e^{-i\Lambda x}\mathbf{z}_2) \end{pmatrix}$$

Renormalization

- Now, assume that $(\lambda_1, \dots, \lambda_d)$ satisfies the GBD condition.

Approximate solution

$$\hat{\mathbf{Z}}(x) = \mathbf{W}(x) + \epsilon N_1(\mathbf{W}(x), x) := \Phi_1(\mathbf{W}(x), x),$$

with $N_1(\mathbf{W}, x) = \int_0^x (\mathbf{g}(\mathbf{W}, s) - R_1(\mathbf{W})) ds$, and $\mathbf{W}(x) = (\mathbf{w}_1^T(x), \mathbf{w}_2^T(x))^T$ is the solution of the RG equation:

$$\begin{cases} \mathbf{W}' = \epsilon R_1(\mathbf{W}), \\ \mathbf{W}(0) = \begin{pmatrix} \frac{\eta_1 - i\Lambda^{-1}\eta_2}{2} \\ \frac{-i\eta_1 + \Lambda^{-1}\eta_2}{2} \end{pmatrix}, \end{cases} \quad (23)$$

$$R_1(\mathbf{W}) = \frac{iA}{2} \begin{pmatrix} \Lambda^{-1}\mathbf{w}_1 \\ -\mathbf{w}_2 \end{pmatrix} + \frac{1}{2} \lim_{x \rightarrow \infty} \frac{1}{x} \int_0^x \begin{pmatrix} i\Lambda^{-1} e^{-i\Lambda s} \mathbf{f}(e^{i\Lambda s} \mathbf{w}_1 + i\Lambda e^{-i\Lambda s} \mathbf{w}_2) \\ -\Lambda^{-1} e^{i\Lambda s} \mathbf{f}(e^{i\Lambda s} \mathbf{w}_1 + i\Lambda e^{-i\Lambda s} \mathbf{w}_2) \end{pmatrix} ds.$$



Approximate results

$$\begin{pmatrix} \hat{\mathbf{y}}_1(t) \\ \hat{\mathbf{y}}_2(t) \end{pmatrix} = \begin{pmatrix} E & i\Lambda \\ iE & \Lambda \end{pmatrix} \begin{pmatrix} e^{i\Lambda t/\varepsilon^2} & O \\ O & e^{-i\Lambda t/\varepsilon^2} \end{pmatrix} \Phi_1(\mathbf{V}(t), t/\varepsilon^2),$$

where

$$\begin{cases} \dot{\mathbf{V}} = R_1(\mathbf{V}), \\ \mathbf{V}(0) = \begin{pmatrix} \frac{\eta_1 - i\Lambda^{-1}\eta_2}{2} \\ \frac{-i\eta_1 + \Lambda^{-1}\eta_2}{2} \end{pmatrix}. \end{cases} \quad (24)$$

with maximum existence interval is (a, b) with $a < b \in \mathbb{R} \cup \{\pm\infty\}$.

Approximate results

Theorem 3.

Assume vector $(\lambda_1, \dots, \lambda_d)$ satisfies the GBD condition. Let $\mathbf{y}(t)$ be the solution of (18), and $\mathbf{V}(t)$ be the solution of (24). Then, for any closed interval $[T_1, T_2] \subset (a, b)$, there exists a constant $\varepsilon_0 > 0$, such that, for any $0 < \varepsilon < \varepsilon_0$,

$$\|\mathbf{y}(t) - \hat{\mathbf{y}}_1(t)\| < C\varepsilon^2, \quad (25)$$

$$\|\dot{\mathbf{y}}(t) - \frac{\hat{\mathbf{y}}_2(t)}{\varepsilon^2}\| < C, \quad (26)$$

$$\left\| \frac{d^k}{dt^k} \mathbf{y}(t) \right\| \leq \frac{M}{\varepsilon^{2k}}, \quad (27)$$

as long as and $T_1 \leq t \leq T_2$, where C, M are positive constants independent to ε .

Original problems

Corollary 2.

Assume that $B = E$, A is a semi-positive definite matrix, $\eta_1, \eta_2 \in \mathbb{R}^d$, and $\mathbf{f}(\mathbf{y}) = \frac{\partial}{\partial \mathbf{y}} h(|\mathbf{y}|^2)$ with $h(r)$ is a polynomial function in \mathbb{R} . Let $\mathbf{y}(t)$ be the solution of the singular perturbed problem (18), $\mathbf{V}(t)$ be the solution of (24). Then, for any $T > 0$, there exists a constant $\varepsilon_0 > 0$, such that, for any $k \in \mathbb{N}$, $0 < \varepsilon < \varepsilon_0$ and $-T \leq t \leq T$,

$$|\mathbf{y}(t) - \hat{\mathbf{y}}_1(t)| < C\varepsilon^2,$$

$$\left| \dot{\mathbf{y}}(t) - \frac{\hat{\mathbf{y}}_2(t)}{\varepsilon^2} \right| < M,$$

$$|\mathbf{y}^{(k)}(t)| \leq \frac{M}{\varepsilon^{2k}},$$

where M is a positive constant independent to ε .

More discussions

- For $d = 1$, by Corollary 2. we have, in fact, obtained the positive answer to the assumption (3) proposed in [Bao *et al* , Numer. Math. 2012; J. Math. Study. 2014].
- The equation (24) is equivalent to

$$\begin{cases} i\dot{\mathbf{v}}_1 + \frac{1}{2}A\mathbf{v}_1 + \frac{1}{2} \sum_{k=1}^N \left(\sum_{l=1}^k C_k^l C_{k-1}^{l-1} \right) kh_k |\mathbf{v}_1|^{2k-2} \mathbf{v}_1 = 0, \\ \mathbf{v}_1(0) = \frac{\eta_1 - i\eta_2}{2}, \end{cases}$$

which is, in fact, a complex Hamiltonian system.



Other problems

$$\begin{cases} \ddot{\mathbf{y}}(t) + \frac{1}{\varepsilon^2} A \mathbf{y}(t) + \mathbf{f}(\mathbf{y}(t)) = 0, \mathbf{y} \in \mathbb{C}^d, t > 0, \\ \mathbf{y}(0) = \eta_1, \dot{\mathbf{y}}(0) = \frac{\eta_2}{\varepsilon}, \end{cases} \quad (28)$$

$$\begin{cases} \mathbf{i} \dot{\mathbf{y}}(t) + \frac{1}{\varepsilon} A \mathbf{y}(t) + \mathbf{f}(\mathbf{y}(t)) = 0, \mathbf{y} \in \mathbb{C}^d, t > 0, \\ \mathbf{y}(0) = \eta_1, \end{cases} \quad (29)$$

- ε : small real parameter $|\varepsilon| \ll 1$;
- A : nonnegative (positive) definite matrix;



Conclusions.



- 1 Gave a systematic introduction to the singular perturbation renormalization group theory under more general conditions.
- 2 Presented a new sufficient condition for quasi-periodic cases to the singular perturbation RG theory.
- 3 Obtained a complete discussion to the highly oscillatory problem (18), and positive answer to the assumption (3) proposed in [Bao *et al* , Numer. Math. 2012; J. Math. Study. 2014].



Thank you!

EMAIL: lwlei@jlu.edu.cn

